

An Eigenfunction Expansion for a Quadratic Pencil of a Schrödinger Operator with Spectral Singularities

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In this paper, we consider the operator L generated in $L^2(\mathbf{R}_+)$ by the differential expression

$$\ell(y) = -y'' + [q(x) + 2\lambda p(x) - \lambda^2]y, \quad x \in \mathbf{R}_+ = [0, \infty),$$

and the boundary condition $y(0) = 0$, where p and q are complex-valued functions and p is continuously differentiable on \mathbf{R}_+ . We derive a two-fold spectral expansion of L (in the sense of Keldysh, 1951, *Soviet Math. Dokl.* **77**, 11–14 [1971, *Russian Math. Survey* **26**, 15–44 (Engl. transl.)]) in terms of the principal functions under the conditions

$$\lim_{x \rightarrow \infty} p(x) = 0, \quad \sup_{x \in \mathbf{R}_+} \{e^{\varepsilon x} [|q(x)| + |p'(x)|]\} < \infty, \quad \varepsilon > 0,$$

taking into account the spectral singularities. Also we investigate the convergence of the spectral expansion. © 1999 Academic Press

1. INTRODUCTION

The spectral analysis of the non-selfadjoint abstract operators with purely discrete spectrum has been considered by Keldysh [10]. He studied the spectrum and principal functions (eigenfunctions and associated functions) of operators involving a polynomial dependent on the spectral parameter, and also showed the completeness of the principal functions of these operators in Hilbert space.

The spectral analysis of a non-selfadjoint differential operators with continuous and discrete spectrum was investigated by Naimark [24]. He

showed the existence of spectral singularities in the continuous spectrum of the non-selfadjoint differential operator L_0 , generated in $L^2(\mathbf{R}_+)$, by the differential expression

$$\ell_0(y) = -y'' + q(x)y, \quad x \in \mathbf{R}_+ = [0, \infty), \quad (1.1)$$

with the boundary condition $y'(0) - hy(0) = 0$, where q is a complex valued function, and $h \in \mathbf{C}$. Also, if $e^{\varepsilon x}q(x) \in L^1(\mathbf{R}_+)$ for some $\varepsilon > 0$, then L_0 has a finite number of eigenvalues and spectral singularities with finite multiplicities. Moreover, he derived the spectral expansion in some particular cases.

Another approach for the discussion of the spectral analysis of L_0 was given by Marchenko [22]. Let E denote the set of all even entire functions of exponential type which are integrable over the real axis, and let E' denote the dual of E . We define

$$\varphi(f_i, \lambda) = \int_0^\infty f_i(x) \varphi(x, \lambda) dx, \quad i = 1, 2,$$

for any finite $f_1, f_2 \in L^2(\mathbf{R}_+)$, where $\varphi(x, \lambda)$ is the solution of $\ell_0(y) = \lambda^2 y$, subject to the initial conditions $\varphi(0, \lambda) = 1$, $\varphi_x(0, \lambda) = h$. In [22] Marchenko showed that

$$\varphi(f_1, \lambda), \varphi(f_2, \lambda) \in E,$$

and there exists a functional $T \in E'$ such that

$$\int_0^\infty f_1(x) f_2(x) dx = T[\varphi(f_1, \lambda) \cdot \varphi(f_2, \lambda)]. \quad (1.2)$$

T is the generalized spectral function of L_0 . (1.2) is a generalization of the well-known Parseval equality for the singular selfadjoint differential operators, and is called Marchenko–Parseval equality.

The results of Naimark were extended to differential operators on the entire real axis by Kemp [11], and to the three dimensional Schrödinger operators by Gasymov [4].

Pavlov [26] established the dependence of the structure of the spectral singularities of L_0 on the behaviour of the potential function at infinity. He first gave the Marchenko–Parseval equality in the form of

$$\int_0^\infty f_1(x) f_2(x) dx = \frac{1}{2\pi i} \int_\gamma m(\lambda) \varphi(f_1, \lambda) \cdot \varphi(f_2, \lambda) d\lambda, \quad (1.3)$$

where γ is the contour enclosing the spectrum of L_0 and $m(\lambda)$ is the Weyl–Titchmarsh function [29]. Then he obtained the spectral expansion of L_0

in terms of principal functions, using (1.3) and the analytical properties of $m(\lambda)$.

In Gasymov's paper [3] the results of Naimark and Marchenko were extended to the case when the potential q of L_0 is equal to $\ell(\ell+1)x^{-2} + V(x)$, where V is a complex valued function, and

$$\int_0^\infty e^{\varepsilon x} |V(x)| dx < \infty,$$

for some $\varepsilon > 0$.

One very important step in the spectral analysis of L_0 was taken by Lyance [18]. He showed that the spectral singularities play an important role in the spectral analysis of L_0 . He also investigated the effect of the spectral singularities in the spectral expansion.

The eigenfunction expansion of the non-selfadjoint operator, generated in $L^2(\mathbf{R}_+)$ by (1.1) and the boundary condition

$$\int_0^\infty K(x) y(x) dx + \alpha y'(0) - \beta y(0) = 0,$$

where $K \in L^2(\mathbf{R}_+)$ is a complex valued function, and $\alpha, \beta \in \mathbf{C}$, with $|\alpha| + |\beta| \neq 0$, was investigated in detail by Krall [12–16].

Note that the principal functions corresponding to the spectral singularities are not the elements of $L^2(\mathbf{R}_+)$ (i.e., these functions increase polynomially as $x \rightarrow \infty$). Also, the spectral singularities belong to the continuous spectrum and are the poles of the resolvent's kernel; but they are not eigenvalues. However, spectral singularities play a certain critical role in the spectral expansion. Their existence is accompanied by specific phenomenon which are new in the sense that they do not occur either in the spectral theory of selfadjoint or normal operators.

Lyance [18] studied the role of spectral singularities in the spectral expansion with respect to principal functions for the Sturm–Liouville operator L_0 by use of Fourier L_0 -transforms. In order to utilize this technique, the potential of L_0 should satisfy Naimark condition

$$\int_0^\infty e^{\varepsilon x} |q(x)| dx < \infty, \quad \varepsilon > 0.$$

The Laurent expansion of the resolvents of the abstract non-selfadjoint operators in neighborhood of spectral singularities was investigated by Gasymov–Maksudov [5] and Maksudov–Allakhverdiev [20]. They also studied the effect of spectral singularities in the spectral analysis of these operators.

The spectral analysis of some classes of dissipative operators with spectral singularities was considered by Pavlov [27] using the theory of functional models [23] and scattering theory [17].

Some problems of spectral theory of differential and some others types of operators with spectral singularities were also studied in [7, 25, 28, 30].

Let us consider the quadratic pencil of the Schrödinger operator L generated in $L^2(\mathbf{R}_+)$ by the differential expression

$$\ell(y) = -y'' + [q(x) + 2\lambda p(x) - \lambda^2] y, \quad x \in \mathbf{R}_+,$$

and the boundary condition $y(0) = 0$, where p and q are complex valued functions and p is continuously differentiable on \mathbf{R}_+ .

Let \tilde{L} denote the operator generated in $L^2(\mathbf{R}_+)$ by the differential expression

$$\tilde{\ell}(y) = -y'' - [\lambda - p(x)]^2 y, \quad x \in \mathbf{R}_+,$$

and the boundary condition $y(0) = 0$. Note that in relativistic quantum mechanics the equation

$$y'' + [\lambda - p(x)]^2 y = 0, \quad x \in \mathbf{R}_+,$$

is called the Klein–Gordon s -wave equation for a particle of zero mass with static potential p [6]. The operators L_0 and \tilde{L} are particular case of L .

Some problems of spectral theory of L and \tilde{L} were studied by Degasperis [2], Jaulent–Jean [8, 9], Maksudov [19] and Maksudov–Guseinov [21].

In this paper, we derive a two-fold spectral expansion of L in terms of the principal functions, under the conditions

$$\lim_{x \rightarrow \infty} p(x) = 0, \quad \sup_{x \in \mathbf{R}_+} \{e^{\varepsilon x} [|q(x)| + |p'(x)|]\} < \infty, \quad \varepsilon > 0, \quad (1.4)$$

taking into account the spectral singularities. Moreover the convergence of the spectral expansion is investigated. Note that we investigate the effect of the spectral singularities in the spectral expansion of L by the regularization of divergent integrals using summability factors.

2. SPECIAL SOLUTIONS OF $\ell(y) = 0$.

Related with the operator L we will consider the boundary value problem

$$-y'' + [q(x) + 2\lambda p(x) - \lambda^2] y = 0, \quad x \in \mathbf{R}_+, \quad (2.1)$$

$$y(0) = 0, \quad (2.2)$$

assuming that the condition (1.4) holds.

Under the condition (1.4) the equation (2.1) has the solutions

$$e^+(x, \lambda) = e^{i w(x) + i \lambda x} + \int_x^\infty A^+(x, t) e^{i \lambda t} dt, \quad (2.3)$$

and

$$e^-(x, \lambda) = e^{-i w(x) - i \lambda x} + \int_x^\infty A^-(x, t) e^{-i \lambda t} dt, \quad (2.4)$$

for $\lambda \in \bar{\mathbf{C}}_+ = \{\lambda: \lambda \in \mathbf{C}, \operatorname{Im} \lambda \geq 0\}$, and $\lambda \in \bar{\mathbf{C}}_- = \{\lambda: \lambda \in \mathbf{C}, \operatorname{Im} \lambda \leq 0\}$, respectively, where $w(x) = \int_x^\infty p(t) dt$, and the kernels $A^\pm(x, t)$ are expressed in terms of p and q , and

$$|A^\pm(x, t)| \leq c \xi \left(\frac{x+t}{2} \right) \exp\{\zeta(x)\}, \quad (2.5)$$

where

$$\xi(x) = \int_x^\infty [|q(t)| + |p'(t)|] dt, \quad \zeta(x) = \int_x^\infty [t |q(t)| + 2 |p(t)|] dt, \quad (2.6)$$

and $c > 0$ is a constant. Therefore $e^+(x, \lambda)$ and $e^-(x, \lambda)$ are analytic with respect to λ in $\mathbf{C}_+ = \{\lambda: \lambda \in \mathbf{C}, \operatorname{Im} \lambda > 0\}$ and $\mathbf{C}_- = \{\lambda: \lambda \in \mathbf{C}, \operatorname{Im} \lambda < 0\}$, respectively, and continuous up to the real axis. $e^\pm(x, \lambda)$ also satisfy

$$e^\pm(x, \lambda) = e^{\pm i[w(x) + \lambda x]} + O\left(\frac{e^{\mp x \operatorname{Im} \lambda}}{|\lambda|}\right), \quad \lambda \in \bar{\mathbf{C}}_\pm, \quad |\lambda| \rightarrow \infty, \quad (2.7)$$

$$e_x^\pm(x, \lambda) = \pm i[\lambda - p(x)] e^{\pm i[w(x) + \lambda x]} + O(1), \quad \lambda \in \bar{\mathbf{C}}_\pm, \quad |\lambda| \rightarrow \infty. \quad (2.8)$$

Let $\varphi(x, \lambda)$ be the solution of (2.1) subject to the initial conditions $\varphi(0, \lambda) = 0$, $\varphi_x(0, \lambda) = 1$. $\varphi(x, \lambda)$ is an entire function of λ and satisfies

$$\varphi(x, \lambda) = \frac{1}{\lambda} \sin \left\{ \lambda x - \int_0^\infty p(t) dt \right\} + O\left(\frac{e^{x \operatorname{Im} \lambda}}{|\lambda|^2}\right), \quad |\lambda| \rightarrow \infty. \quad (2.9)$$

The results stated above were obtained by Jaulent–Jean [8].

3. THE SPECTRUM OF L

By $\sigma_c(L)$, $\sigma_d(L)$ and $\sigma_{ss}(L)$ we denote the continuous spectrum, the eigenvalues and the spectral singularities of L , respectively. We have previously shown [1] that

$$\sigma_c(L) = (-\infty, \infty),$$

$$\sigma_d(L) = \{\lambda: \lambda \in \mathbf{C}_+, e^+(\lambda) = 0\} \cup \{\lambda: \lambda \in \mathbf{C}_-, e^-(\lambda) = 0\},$$

$$\sigma_{ss}(L) = \{\lambda: \lambda \in \mathbf{R}^*, e^+(\lambda) = 0\} \cup \{\lambda: \lambda \in \mathbf{R}^*, e^-(\lambda) = 0\},$$

where

$$e^{\pm}(\lambda) = e^{\pm}(0, \lambda), \quad \mathbf{R}^* = \mathbf{R} \setminus \{0\}.$$

Under the condition (1.4), we know that L has a finite number of eigenvalues and spectral singularities, and each of them is of finite multiplicity [1]. Let $\lambda_1^+, \dots, \lambda_j^+$ and $\lambda_1^-, \dots, \lambda_k^-$ denote the zeros of the functions e^+ in \mathbf{C}_+ and e^- in \mathbf{C}_- (which are the eigenvalues of L) with multiplicities m_1^+, \dots, m_j^+ and m_1^-, \dots, m_k^- , respectively. Similarly, let $\lambda_1, \dots, \lambda_v$ and $\lambda_{v+1}, \dots, \lambda_\ell$ be zeros of e^+ and e^- in \mathbf{R}^* (the spectral singularities of L) with multiplicities n_1, \dots, n_v and n_{v+1}, \dots, n_ℓ , respectively.

We will also need the Hilbert spaces

$$H_n = \left\{ f: \int_0^\infty (1+x)^{2n} |f(x)|^2 dx < \infty \right\}, \quad n = 1, 2, \dots,$$

$$H_{-n} = \left\{ g: \int_0^\infty (1+x)^{-2n} |g(x)|^2 dx < \infty \right\}, \quad n = 1, 2, \dots,$$

with

$$\|f\|_n^2 = \int_0^\infty (1+x)^{2n} |f(x)|^2 dx, \quad \|g\|_{-n}^2 = \int_0^\infty (1+x)^{-2n} |g(x)|^2 dx,$$

respectively. It is clear that,

$$H_{(n+1)} \subsetneq H_n \subsetneq L^2(\mathbf{R}_+) \subsetneq H_{-n} \subsetneq H_{-(n+1)}, \quad n = 1, 2, \dots,$$

and H_{-n} is isomorphic to the dual of H_n : $H'_n \sim H_{-n}$. In fact, for every functional $F \in H'_n$, there is a function f^* belonging to H_{-n} such that

$$F(f) = \int_0^\infty f(x) f^*(x) dx$$

for all $f \in H_n$.

We have previously shown that [1]:

$$\left\{ \left\{ \frac{\partial^n}{\partial \lambda^n} \varphi(\cdot, \lambda) \right\}_{\lambda=\lambda_i^+} \right\} \in L^2(\mathbf{R}_+), \quad n = 0, 1, \dots, m_i^+ - 1, \quad i = 1, \dots, j, \left\{ \left\{ \frac{\partial^n}{\partial \lambda^n} \varphi(\cdot, \lambda) \right\}_{\lambda=\lambda_i^-} \right\} \in L^2(\mathbf{R}_+), \quad n = 0, 1, \dots, m_i^- - 1, \quad i = 1, \dots, k, \quad (3.1)$$

$$\left\{ \left\{ \frac{\partial^n}{\partial \lambda^n} \varphi(\cdot, \lambda) \right\}_{\lambda=\lambda_i} \right\} \in H_{-(n+1)}, \quad n = 0, 1, \dots, n_i - 1, \quad i = 1, \dots, v, v+1, \dots, \ell. \quad (3.2)$$

LEMMA 3.1. For $\lambda \in \mathbf{R}^*$,

$$\frac{\partial^n}{\partial \lambda^n} \varphi(\cdot, \lambda) \in H_{-(n+1)}. \quad (3.3)$$

Proof. It is known that, for $\lambda \in \mathbf{R}^*$,

$$2i\lambda\varphi(x, \lambda) = e^-(\lambda) e^+(x, \lambda) - e^+(\lambda) e^-(x, \lambda). \quad (3.4)$$

Then the proof of the lemma is the direct consequence of (2.3), (2.4), and (3.4). ■

Let

$$G(x, t; \lambda) = \begin{cases} G^+(x, t; \lambda), & \lambda \in \mathbf{C}_+, \\ G^-(x, t; \lambda), & \lambda \in \mathbf{C}_-, \end{cases} \quad (3.5)$$

be the Green function of L , where

$$G^\pm(x, t; \lambda) = \frac{1}{e^\pm(\lambda)} \begin{cases} e^\pm(x, \lambda) \varphi(t, \lambda), & 0 \leq t < x, \\ e^\pm(t, \lambda) \varphi(x, \lambda), & x \leq t < \infty. \end{cases} \quad (3.6)$$

4. THE SPECTRAL EXPANSION

Let $C_0^\infty(\mathbf{R}_+)$ denote the set of infinitely differentiable functions in \mathbf{R}_+ with compact support. Then

$$f(x) = \int_0^\infty G(x, t; \lambda) [-f''(t) + q(t)f(t) + 2\lambda p(t)f(t) - \lambda^2 f(t)] dt,$$

for each $f \in C_0^\infty(\mathbf{R}_+)$. Hence we have

$$\frac{f(x)}{\lambda} = \frac{1}{\lambda} \int_0^\infty G(x, t; \lambda) \Theta(t) dt + 2 \int_0^\infty G(x, t; \lambda) p(t) f(t) dt - \lambda D(x, \lambda), \quad (4.1)$$

where

$$\Theta(t) = -f''(t) + q(t)f(t), \quad D(x, \lambda) = \int_0^\infty G(x, t; \lambda) f(t) dt.$$

Let γ_r denote the disc with center at the origin having radius r ; let $\partial\gamma_r$ be the boundary of γ_r . r will be chosen so that all eigenvalues and spectral singularities of L are in γ_r . $P_{r\eta}$ denotes the part of γ_r lying in the strip $|\operatorname{Im} \lambda| \leq \eta$ and $\gamma_{r\eta} = \gamma_{r\eta}^+ \cup \gamma_{r\eta}^-$, where $\gamma_{r\eta}^+$ and $\gamma_{r\eta}^-$ are the parts of $\gamma_r - P_{r\eta}$ in the upper and lower half-planes, respectively (see Fig. 1). We choose η so small that $P_{r\eta}$ does not contain any eigenvalues of L .

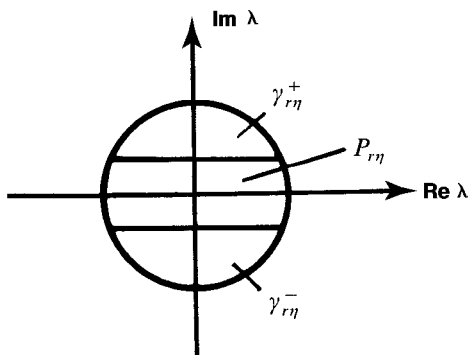


FIGURE 1

As it is seen in Fig.1,

$$\partial\gamma_r = \partial\gamma_{r\eta} \cup \partial P_{r\eta}. \quad (4.2)$$

From (4.1) we obtain

$$\begin{aligned} f(x) = & \frac{1}{2\pi i} \int_{\partial\gamma_r} \left\{ \frac{1}{\lambda} \int_0^\infty G(x, t; \lambda) \Theta(t) dt \right\} d\lambda \\ & + \frac{1}{\pi i} \int_{\partial\gamma_r} \left\{ \int_0^\infty G(x, t; \lambda) p(t) f(t) dt \right\} d\lambda \\ & - \frac{1}{2\pi i} \int_{\partial\gamma_r} \lambda D(x, \lambda) d\lambda. \end{aligned} \quad (4.3)$$

Using (2.7), (2.9), (3.5), (3.6) and Jordan's lemma, we see that the first term of the right-hand side of (4.3) vanishes as $r \rightarrow \infty$. The same result holds for the second term. This can be obtained from (2.7)–(2.9) utilizing integration by parts. Then considering (4.2) we find

$$f(x) = - \lim_{\substack{r \rightarrow \infty \\ \eta \rightarrow 0}} \frac{1}{2\pi i} \int_{\partial\gamma_{r\eta}} \lambda D(x, \lambda) d\lambda - \lim_{\substack{r \rightarrow \infty \\ \eta \rightarrow 0}} \frac{1}{2\pi i} \int_{\partial P_{r\eta}} \lambda D(x, \lambda) d\lambda. \quad (4.4)$$

Therefore

$$\frac{1}{2\pi i} \int_{\partial\gamma_{r\eta}} \lambda D(x, \lambda) d\lambda = \sum_{i=1}^j \operatorname{Res}_{\lambda=\lambda_i^+} [\lambda D^+(x, \lambda)] + \sum_{i=1}^k \operatorname{Res}_{\lambda=\lambda_i^-} [\lambda D^-(x, \lambda)],$$

where

$$D^\pm(x, \lambda) = \int_0^\infty G^\pm(x, t; \lambda) f(t) dt.$$

Let Γ_+ be the contour which isolates the real zeros of e^+ by semicircles with centers at λ_i , $i = 1, \dots, v$, having the same radius δ_0 in the upper half-plane. Similarly Γ_- will denote the corresponding contour for the real zeros of e^- in the lower half-plane. The radius of semicircles being chosen so small that their diameters are mutually disjoint and do not contain the point $\lambda = 0$ (see Fig. 2).

As it is easily seen from Fig. 1, we find

$$\lim_{\substack{r \rightarrow \infty \\ \eta \rightarrow 0}} \frac{1}{2\pi i} \int_{P_{\gamma\eta}} \lambda D(x, \lambda) d\lambda = \frac{1}{2\pi i} \int_{\Gamma_-} \lambda D^-(x, \lambda) d\lambda - \frac{1}{2\pi i} \int_{\Gamma_+} \lambda D^+(x, \lambda) d\lambda$$

So (4.4) can be written as

$$\begin{aligned} f(x) = & - \sum_{i=1}^j \operatorname{Res}_{\lambda=\lambda_i^+} [\lambda D^+(x, \lambda)] - \sum_{i=1}^k \operatorname{Res}_{\lambda=\lambda_i^-} [\lambda D^-(x, \lambda)] \\ & + \frac{1}{2\pi i} \int_{\Gamma_+} \lambda D^+(x, \lambda) d\lambda - \frac{1}{2\pi i} \int_{\Gamma_-} \lambda D^-(x, \lambda) d\lambda. \end{aligned} \quad (4.5)$$

LEMMA 4.1. For every $f \in C_0^\infty(\mathbf{R}_+)$

$$\begin{aligned} f(x) = & \sum_{i=1}^j \left\{ \left(\frac{\partial}{\partial \lambda} \right)^{m_i^+ - 1} [a_i^+(\lambda) \varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_i^+} \\ & + \sum_{i=1}^k \left\{ \left(\frac{\partial}{\partial \lambda} \right)^{m_i^- - 1} [a_i^-(\lambda) \varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_i^-} \\ & + \frac{1}{2\pi i} \int_{\Gamma_+} \frac{\lambda e_x^+(0, \lambda)}{e^+(\lambda)} \varphi(x, \lambda) \varphi(f, \lambda) d\lambda \\ & - \frac{1}{2\pi i} \int_{\Gamma_-} \frac{\lambda e_x^-(0, \lambda)}{e^-(\lambda)} \varphi(x, \lambda) \varphi(f, \lambda) d\lambda, \end{aligned} \quad (4.6)$$

$$\begin{aligned} 0 = & \sum_{i=1}^j \left\{ \left(\frac{\partial}{\partial \lambda} \right)^{m_i^+ - 1} [b_i^+(\lambda) \varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_i^+} \\ & + \sum_{i=1}^k \left\{ \left(\frac{\partial}{\partial \lambda} \right)^{m_i^- - 1} [b_i^-(\lambda) \varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_i^-} \\ & + \frac{1}{2\pi i} \int_{\Gamma_+} \frac{e_x^+(0, \lambda)}{e^+(\lambda)} \varphi(x, \lambda) \varphi(f, \lambda) d\lambda \\ & - \frac{1}{2\pi i} \int_{\Gamma_-} \frac{e_x^-(0, \lambda)}{e^-(\lambda)} \varphi(x, \lambda) \varphi(f, \lambda) d\lambda, \end{aligned} \quad (4.7)$$

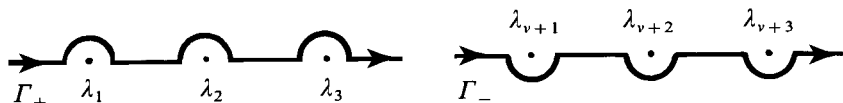


FIGURE 2

where

$$\left. \begin{aligned} a_i^+(\lambda) &= -\frac{\lambda(\lambda - \lambda_i^+)^{m_i^+} e_x^+(0, \lambda)}{(m_i^+ - 1)! e^+(\lambda)}, & i = 1, \dots, j, \\ a_i^-(\lambda) &= -\frac{\lambda(\lambda - \lambda_i^-)^{m_i^-} e_x^-(0, \lambda)}{(m_i^- - 1)! e^-(\lambda)}, & i = 1, \dots, k, \end{aligned} \right\} \quad (4.8)$$

$$\left. \begin{aligned} b_i^+(\lambda) &= -\frac{(\lambda - \lambda_i^+)^{m_i^+} e_x^+(0, \lambda)}{(m_i^+ - 1)! e^+(\lambda)}, & i = 1, \dots, j, \\ b_i^-(\lambda) &= -\frac{(\lambda - \lambda_i^-)^{m_i^-} e_x^-(0, \lambda)}{(m_i^- - 1)! e^-(\lambda)}, & i = 1, \dots, k, \end{aligned} \right\} \quad (4.9)$$

and

$$\varphi(f, \lambda) = \int_0^\infty f(t) \varphi(t, \lambda) dt.$$

Proof. Let $\psi(x, \lambda)$ be the solution of (2.1) subject to the initial conditions $\psi(0, \lambda) = 1$, $\psi_x(0, \lambda) = 0$. Then

$$G^\pm(x, t; \lambda) = \frac{e_x^\pm(0, \lambda)}{e^\pm(\lambda)} \varphi(x, \lambda) \varphi(t, \lambda) + a(x, t; \lambda), \quad (4.10)$$

where

$$a(x, t; \lambda) = \begin{cases} \psi(x, \lambda) \varphi(t, \lambda), & 0 \leq t < x, \\ \psi(t, \lambda) \varphi(x, \lambda), & x \leq t < \infty, \end{cases}$$

and $a(x, t; \lambda)$ is an entire function of λ . From (4.5) and (4.10) we obtain (4.6). Writing (4.1) as

$$\frac{f(x)}{\lambda^2} = \frac{1}{\lambda^2} \int_0^\infty G(x, t; \lambda) \Theta(t) dt + \frac{2}{\lambda} \int_0^\infty G(x, t; \lambda) p(t) f(t) dt - D(x, \lambda),$$

and repeating the calculation as we have done for (4.1), we have (4.7). ■

Since the contours Γ_+ and Γ_- in (4.6) and (4.7) do not coincide with the continuous spectrum of L , these formulae contains “non-spectral”

objects. The purpose of this article is to transform (4.6) and (4.7) into a two-fold spectral expansion with respect to the principal functions of L .

LEMMA 4.2. *For any $f \in C_0^\infty(\mathbf{R}_+)$ there exists a constant $c > 0$ so that*

$$\int_{-\infty}^{\infty} |\lambda \varphi(f, \lambda)|^2 d\lambda \leq c \int_0^{\infty} |f(x)|^2 dx. \quad (4.11)$$

Proof. From (3.4) we get

$$2i\lambda \varphi(f, \lambda) = e^-(\lambda) e^+(f, \lambda) - e^+(\lambda) e^-(f, \lambda), \quad \lambda \in \mathbf{R}^*, \quad (4.12)$$

where

$$e^\pm(f, \lambda) = \int_0^{\infty} f(x) e^\pm(x, \lambda) dx.$$

Using (2.3) we have

$$e^+(f, \lambda) = \int_0^{\infty} [(\mathbf{A}_0^+ + \mathbf{A}^+) f(t)] e^{i\lambda t} dt, \quad (4.13)$$

in which the operators \mathbf{A}_0^+ and \mathbf{A}^+ are defined by

$$\mathbf{A}_0^+ f(t) = e^{i\omega(t)} f(t), \quad \mathbf{A}^+ f(t) = \int_0^t A^+(x, t) f(x) dx.$$

From (1.4) and (2.5) we find

$$|A^+(x, t)| \leq c_0 \exp \left\{ -\frac{\varepsilon}{4} (x + t) \right\},$$

where $c_0 > 0$ is a constant. Hence \mathbf{A}^+ is a compact operator in $L^2(\mathbf{R}_+)$. Thus $(\mathbf{A}_0^+ + \mathbf{A}^+)$ is continuous and one-to-one on $L^2(\mathbf{R}_+)$. Using the Parseval's equality for the Fourier transforms and (4.13) we get

$$\int_{-\infty}^{\infty} |e^+(f, \lambda)|^2 d\lambda \leq c_1 \int_0^{\infty} |f(x)|^2 dx, \quad (4.14)$$

where $c_1 > 0$ is a constant. In a similar way, we also find

$$\int_{-\infty}^{\infty} |e^-(f, \lambda)|^2 d\lambda \leq c_2 \int_0^{\infty} |f(x)|^2 dx. \quad (4.15)$$

The proof of the lemma is completed by (2.7), (4.12), (4.14) and (4.15). ■

By the preceding lemma for every function $f \in L^2(\mathbf{R}_+)$ the limit

$$\varphi(f, \lambda) = \lim_{N \rightarrow \infty} \int_0^N f(x) \varphi(x, \lambda) dx$$

exists in the sense of convergence in the mean square, relative to the measure $\lambda^2 d\lambda$ on the real axis; that is,

$$\lim_{N \rightarrow \infty} \int_{-\infty}^{\infty} \left| \varphi(f, \lambda) - \int_0^N f(x) \varphi(x, \lambda) dx \right|^2 \lambda^2 d\lambda = 0. \quad (4.16)$$

Since $C_0^\infty(\mathbf{R}_+)$ is dense in $L^2(\mathbf{R}_+)$, the estimate (4.11) may be extended onto $L^2(\mathbf{R}_+)$ for any $f \in L^2(\mathbf{R}_+)$ as

$$\int_{-\infty}^{\infty} |\lambda \varphi(f, \lambda)|^2 d\lambda \leq c \int_0^{\infty} |f(x)|^2 dx, \quad (4.17)$$

where $\varphi(f, \lambda)$ must be understood in the sense of (4.16). We shall need a generalization of this estimate.

LEMMA 4.3. *If $f \in H_m$, then $\varphi(f, \lambda)$ has a derivative of order $(m-1)$ which is absolutely continuous of every finite subinterval of the real axis and satisfies*

$$\int_{-\infty}^{\infty} \left| \left(\frac{d}{d\lambda} \right)^n [\lambda \varphi(f, \lambda)] \right|^2 d\lambda \leq c_n \int_0^{\infty} (1+x)^{2n} |f(x)|^2 dx, \quad n = 1, \dots, m, \quad (4.18)$$

where $c_n > 0$ are constants, $n = 1, \dots, m$.

The proof is similar to that of Lemma 4.2.

In order to transform (4.6) and (4.7) into the spectral expansion of L , we have to reform the integrals over Γ_+ and Γ_- onto the real axis (i.e., to the continuous spectrum of L).

Since the spectral singularities of L are the real zeros of e^+ and e^- , the integrals over the real axis are divergent in the norm of $L^2(\mathbf{R}_+)$. Now we will investigate the convergence of these integrals in a norm which is weaker than the norm of $L^2(\mathbf{R}_+)$. For this purpose we will use the technique of the regularization of divergent integrals. So we define the following summability factors

$$F_{\alpha\beta}^+(\lambda) = \begin{cases} \frac{(\lambda - \lambda_\alpha)^\beta}{\beta!}, & |\lambda - \lambda_\alpha| < \delta, \quad \alpha = 1, \dots, v, \\ 0, & |\lambda - \lambda_\alpha| \geq \delta, \quad \alpha = 1, \dots, v. \end{cases} \quad (4.19)$$

$$F_{\alpha\beta}^-(\lambda) = \begin{cases} \frac{(\lambda - \lambda_\alpha)^\beta}{\beta!}, & |\lambda - \lambda_\alpha| < \delta, \quad \alpha = v+1, \dots, \ell, \\ 0, & |\lambda - \lambda_\alpha| \geq \delta, \quad \alpha = v+1, \dots, \ell, \end{cases} \quad (4.20)$$

with $\delta > \delta_0$. We can choose $\delta > 0$ so small that the δ -neighborhoods of λ_α , $\alpha = 1, \dots, v, v+1, \dots, \ell$ have no common points and do not contain the point $\lambda = 0$. Define the functionals

$$F^+\{g_1(\lambda)\} = g_1(\lambda) - \sum_{\alpha=1}^v \sum_{\beta=0}^{n_\alpha-1} \left\{ \left(\frac{d}{d\lambda} \right)^\beta g_1(\lambda) \right\}_{\lambda=\lambda_\alpha} F_{\alpha\beta}^+(\lambda), \quad (4.21)$$

$$F^-\{g_2(\lambda)\} = g_2(\lambda) - \sum_{\alpha=v+1}^\ell \sum_{\beta=0}^{n_\alpha-1} \left\{ \left(\frac{d}{d\lambda} \right)^\beta g_2(\lambda) \right\}_{\lambda=\lambda_\alpha} F_{\alpha\beta}^-(\lambda), \quad (4.22)$$

where g_1 and g_2 are chosen so that the right-hand side of the above formulae are meaningful. It is evident from (4.19)–(4.20) that $\lambda_1, \dots, \lambda_v$ and $\lambda_{v+1}, \dots, \lambda_\ell$ are the roots of $F^+\{g_1(\lambda)\} = 0$ and $F^-\{g_2(\lambda)\} = 0$ at least of orders n_1, \dots, n_v and n_{v+1}, \dots, n_ℓ , respectively. In the following we will use the operators

$$P^+f(x) = \frac{1}{2\pi i} \int_{\Gamma_+} \frac{\lambda e_x^+(0, \lambda)}{e^+(\lambda)} \varphi(x, \lambda) \varphi(f, \lambda) d\lambda, \quad (4.23)$$

and

$$P^-f(x) = \frac{1}{2\pi i} \int_{\Gamma_-} \frac{\lambda e_x^-(0, \lambda)}{e^-(\lambda)} \varphi(x, \lambda) \varphi(f, \lambda) d\lambda, \quad (4.24)$$

$$\begin{aligned} I^+f(x) &= \frac{1}{2\pi i} \sum_{\alpha=1}^v \sum_{\beta=0}^{n_\alpha-1} \left\{ \left(\frac{\partial}{\partial \lambda} \right)^\beta [\varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_\alpha} \\ &\quad \times \int_{\Gamma_+} \frac{\lambda e_x^+(0, \lambda)}{e^+(\lambda)} F_{\alpha\beta}^+(\lambda) d\lambda \\ &\quad + \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{\lambda e_x^+(0, \lambda)}{e^+(\lambda)} F^+\{\varphi(x, \lambda) \varphi(f, \lambda)\} d\lambda, \end{aligned}$$

and

$$\begin{aligned} I^-f(x) &= \frac{1}{2\pi i} \sum_{\alpha=v+1}^\ell \sum_{\beta=0}^{n_\alpha-1} \left\{ \left(\frac{\partial}{\partial \lambda} \right)^\beta [\varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_\alpha} \\ &\quad \times \int_{\Gamma_-} \frac{\lambda e_x^-(0, \lambda)}{e^-(\lambda)} F_{\alpha\beta}^-(\lambda) d\lambda \\ &\quad + \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{\lambda e_x^-(0, \lambda)}{e^-(\lambda)} F^-\{\varphi(x, \lambda) \varphi(f, \lambda)\} d\lambda. \end{aligned}$$

Since, under the condition (1.4) $e^+(x, \lambda)$ and $e^-(x, \lambda)$ have an analytic continuation to the half-planes $\operatorname{Im} \lambda > -\varepsilon/2$ and $\operatorname{Im} \lambda < \varepsilon/2$, respectively, we get

$$P^\pm f = I^\pm f,$$

for $f \in C_0^\infty(\mathbf{R}_+)$.

LEMMA 4.4. *For each $f \in H_{(n_0+1)}$, there exist a constant $c > 0$ such that*

$$\|I^\pm f\|_{-(n_0+1)} \leq c \|f\|_{(n_0+1)}, \quad (4.25)$$

where $n_0 = \max\{n_1, \dots, n_v, n_{v+1}, \dots, n_\ell\}$.

Define

$$A_\alpha^+ = (\lambda_\alpha - \delta, \lambda_\alpha + \delta), \quad \alpha = 1, \dots, v. \quad (4.26)$$

Then $0 \notin A_\alpha^+$, $\alpha = 1, \dots, v$. Using the integral form of remainder in the Taylor formula, we get

$$F^+ \{ \varphi(x, \lambda) \varphi(f, \lambda) \} = \begin{cases} \varphi(x, \lambda) \varphi(f, \lambda), & \lambda \in A_0^+ \\ \frac{1}{(n_\alpha - 1)!} \int_{\lambda_\alpha}^{\lambda} (\lambda - \xi)^{n_\alpha - 1} \left\{ \left(\frac{\partial}{\partial \xi} \right)^{n_\alpha} [\varphi(x, \xi) \varphi(f, \xi)] \right\} d\xi \\ \lambda \in A_\alpha^+, & \alpha = 1, \dots, v, \end{cases} \quad (4.27)$$

where

$$A_0^+ = \mathbf{R} \setminus \left\{ \bigcup_{\alpha=1}^v A_\alpha^+ \right\}.$$

If we use the notation

$$I_\alpha^+ f(x) = \frac{1}{2\pi i} \int_{A_\alpha^+} \frac{\lambda e_x^+(0, \lambda)}{e^+(\lambda)} F^+ \{ \varphi(x, \lambda) \varphi(f, \lambda) \} d\lambda, \quad \alpha = 0, 1, \dots, v,$$

$$\tilde{I}^+ f(x) = \frac{1}{2\pi i} \sum_{\alpha=1}^v \sum_{\beta=0}^{n_\alpha-1} \left\{ \left(\frac{\partial}{\partial \lambda} \right)^\beta [\varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_\alpha} \\ \times \int_{\Gamma_+} \frac{\lambda e_x^+(0, \lambda)}{e^+(\lambda)} F_{\alpha\beta}^+(\lambda) d\lambda,$$

we obtain

$$I^+ = I_0^+ + I_1^+ + \dots + I_v^+ + \tilde{I}^+ \quad (4.28)$$

from (4.26) and (4.27). We now show that each of the operators $I_0^+, I_1^+, \dots, I_v^+$ and \tilde{I}^+ is continuous from $H_{(n_0+1)}$ into $H_{-(n_0+1)}$. We start with \tilde{I}^+ . From (4.19) we obtain the absolute convergence of

$$\int_{r_+} \frac{\lambda e_x^+(0, \lambda)}{e^+(\lambda)} F_{\alpha\beta}^+(\lambda) d\lambda.$$

Using (3.2) and the isomorphism $H_{-n_0} \sim H'_{n_0}$ we see that \tilde{I}^+ is continuous from H_{n_0} into H_{-n_0} or from $H_{(n_0+1)}$ into $H_{-(n_0+1)}$. Hence there exists a constant $\tilde{c} > 0$ such that

$$\|\tilde{I}^+ f(x)\|_{-(n_0+1)} \leq \tilde{c} \|f\|_{(n_0+1)} \quad (4.29)$$

for any $f \in H_{(n_0+1)}$.

Next we want to show the continuity of I_α^+ , $\alpha = 1, \dots, v$, from $H_{(n_0+1)}$ into $H_{-(n_0+1)}$. From (4.27) we see that

$$\begin{aligned} I_\alpha^+ f(x) &= \frac{1}{2\pi i(n_\alpha - 1)!} \int_{A_\alpha^+} \frac{\lambda e_x^+(0, \lambda)}{e^+(\lambda)} \int_{\lambda_\alpha}^\lambda (\lambda - \xi)^{n_\alpha - 1} \\ &\quad \times \left\{ \left(\frac{\partial}{\partial \xi} \right)^{n_\alpha} [\varphi(x, \xi) \varphi(f, \xi)] \right\} d\xi d\lambda. \end{aligned} \quad (4.30)$$

Interchanging the order of integration, we get

$$\begin{aligned} I_\alpha^+ f(x) &= \frac{1}{2\pi i(n_\alpha - 1)!} \\ &\quad \times \left\{ \int_{\lambda_\alpha}^{\lambda_\alpha + \delta} \int_{\xi}^{\lambda_\alpha + \delta} \left\{ \left(\frac{\partial}{\partial \xi} \right)^{n_\alpha} [\varphi(x, \xi) \varphi(f, \xi)] \right\} \right. \\ &\quad \times (\lambda - \xi)^{n_\alpha - 1} \frac{\lambda e_x^+(0, \lambda)}{e^+(\lambda)} d\lambda d\xi \\ &\quad \left. - \int_{\lambda_\alpha - \delta}^{\lambda_\alpha} \int_{\lambda_\alpha - \delta}^{\xi} \left\{ \left(\frac{\partial}{\partial \xi} \right)^{n_\alpha} [\varphi(x, \xi) \varphi(f, \xi)] \right\} \right. \\ &\quad \left. \times (\lambda - \xi)^{n_\alpha - 1} \frac{\lambda e_x^+(0, \lambda)}{e^+(\lambda)} d\lambda d\xi \right\}. \end{aligned}$$

Since λ_α is a zero of e^+ order n_α , there exists a continuous function e_α^+ such that $e_\alpha^+(\lambda_\alpha) \neq 0$ and $e^+(\lambda) = (\lambda - \lambda_\alpha)^{n_\alpha} e_\alpha^+(\lambda)$. On the other hand,

$$\left| \int_{\xi}^{\lambda_\alpha + \delta} \frac{\lambda (\lambda - \xi)^{n_\alpha - 1} e_x^+(0, \lambda)}{e^+(\lambda)} d\lambda \right| \leq h_\alpha^+(\xi) [\ln \delta - \ln(\xi - \lambda_\alpha)], \quad (4.31)$$

if $\xi > \lambda_\alpha$, and

$$\left| \int_{\lambda_\alpha - \delta}^{\xi} \frac{\lambda(\lambda - \xi)^{n_\alpha - 1} e_x^+(0, \lambda)}{e^+(\lambda)} d\lambda \right| \leq h_\alpha^-(\xi) [\ln(\lambda_\alpha - \xi) - \ln \delta], \quad (4.32)$$

if $\xi < \lambda_\alpha$, where

$$h_\alpha^+(\xi) = \max_{\lambda \in [\xi, \lambda_\alpha + \delta]} \left| \frac{\lambda e_x^+(0, \lambda)}{e^+(\lambda)} \right|, \quad h_\alpha^-(\xi) = \max_{\lambda \in [\lambda_\alpha - \delta, \xi]} \left| \frac{\lambda e_x^+(0, \lambda)}{e^+(\lambda)} \right|.$$

(4.31) and (4.32) show that I_α^+ , $\alpha = 1, \dots, v$, are integral operators with kernels having logarithmic singularities, i.e., weak singularities. (4.30) can be written as

$$I_\alpha^+ f(x) = \int_{A_\alpha^+} \sum_{k=0}^{n_\alpha} b_{k\alpha}(x, \xi) \left\{ \left(\frac{d}{d\xi} \right)^k \varphi(f, \xi) \right\} d\xi.$$

Define

$$B_{k\alpha} = \int_0^\infty \int_{A_\alpha^+} \left| \frac{b_{k\alpha}(x, \xi)}{(1+x)^{n_0+1}} \right|^2 d\xi dx.$$

We see that $B_{k\alpha} < \infty$, by (3.3), (4.31) and (4.32). Since

$$\begin{aligned} \|I_\alpha^+ f\|_{-(n_0+1)}^2 &= \int_0^\infty \left| \frac{I_\alpha^+ f(x)}{(1+x)^{n_0+1}} \right|^2 dx \\ &\leq \sum_{k=0}^{n_\alpha} \int_0^\infty \int_{A_\alpha^+} \left| \frac{b_{k\alpha}(x, \xi)}{(1+x)^{n_0+1}} \right|^2 d\xi dx \int_{A_\alpha^+} \left| \left(\frac{d}{d\xi} \right)^k \varphi(f, \xi) \right|^2 d\xi \\ &= \sum_{k=0}^{n_\alpha} B_{k\alpha} \int_{A_\alpha^+} \left| \left(\frac{d}{d\xi} \right)^k \varphi(f, \xi) \right|^2 d\xi. \end{aligned}$$

Utilizing (4.17) and (4.18) we obtain

$$\|I_\alpha^+ f\|_{-(n_0+1)} \leq c_\alpha \|f\|_{n_0} \leq c_\alpha \|f\|_{(n_0+1)}, \quad \alpha = 1, \dots, v, \quad (4.33)$$

where c_α are constants.

Lastly we consider the operator I_0^+ which is defined by

$$I_0^+ f(x) = \frac{1}{2\pi i} \int_{-\infty}^\infty \chi_0^+(\lambda) \frac{\lambda e_x^+(0, \lambda)}{e^+(\lambda)} \varphi(x, \lambda) \varphi(f, \lambda) d\lambda, \quad (4.34)$$

where χ_0^+ is the characteristic function of the interval A_0^+ . From (4.34), similar to the proof of Lemma 4.2, we get

$$\int_0^\infty |I_\alpha^+ f(x)|^2 dx \leq c_0 \int_0^\infty |f(x)|^2 dx,$$

where $c_0 > 0$ is a constant. Since

$$H_{(n_0+1)} \subsetneq L^2(\mathbf{R}_+) \subsetneq H_{-(n_0+1)},$$

we find

$$\|I_0^+ f\|_{-(n_0+1)} \leq c_0 \|f\|_{(n_0+1)}. \quad (4.35)$$

From (4.28), (4.29), (4.33), and (4.35) we have

$$\|I^+ f\|_{-(n_0+1)} \leq c \|f\|_{(n_0+1)}.$$

In a similar way it follows that

$$\|I^- f\|_{-(n_0+1)} \leq c \|f\|_{(n_0+1)}. \quad \blacksquare$$

Then for every $f \in H_{(n_0+1)}$,

$$\begin{aligned} I^+ f(x) &= \frac{1}{2\pi i} \sum_{\alpha=1}^v \sum_{\beta=0}^{n_\alpha-1} \left\{ \left(\frac{\partial}{\partial \lambda} \right)^\beta [\varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_\alpha} \\ &\quad \times \int_{\Gamma_+} \frac{\lambda e_x^+(0, \lambda)}{e^+(\lambda)} F_{\alpha\beta}^+(\lambda) d\lambda \\ &\quad + \frac{1}{2\pi i} \int_{-\infty}^\infty \frac{\lambda e_x^+(0, \lambda)}{e^+(\lambda)} F^+ \{ \varphi(x, \lambda) \varphi(f, \lambda) \} d\lambda, \end{aligned} \quad (4.36)$$

and

$$\begin{aligned} I^- f(x) &= \frac{1}{2\pi i} \sum_{\alpha=v+1}^\ell \sum_{\beta=0}^{n_\alpha-1} \left\{ \left(\frac{\partial}{\partial \lambda} \right)^\beta [\varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_\alpha} \\ &\quad \times \int_{\Gamma_+} \frac{\lambda e_x^-(0, \lambda)}{e^-(\lambda)} F_{\alpha\beta}^-(\lambda) d\lambda \\ &\quad + \frac{1}{2\pi i} \int_{-\infty}^\infty \frac{\lambda e_x^-(0, \lambda)}{e^-(\lambda)} F^- \{ \varphi(x, \lambda) \varphi(f, \lambda) \} d\lambda. \end{aligned} \quad (4.37)$$

Let $a_\alpha(\lambda)$ denote any function which is defined and differentiable in a neighborhood of λ_α , and which satisfies the condition

$$\left\{ \left(\frac{d}{d\lambda} \right)^{n_\alpha-1-\beta} a_\alpha(\lambda) \right\}_{\lambda=\lambda_\alpha} = \begin{cases} \frac{1}{2\pi i} \binom{n_\alpha-1}{\beta} \int_{\Gamma_+} \frac{\lambda e_x^+(0, \lambda)}{e^+(\lambda)} F_{\alpha\beta}^+(\lambda) d\lambda, & \alpha = 1, \dots, v \\ -\frac{1}{2\pi i} \binom{n_\alpha-1}{\beta} \int_{\Gamma_-} \frac{\lambda e_x^-(0, \lambda)}{e^-(\lambda)} F_{\alpha\beta}^-(\lambda) d\lambda, & \alpha = v+1, \dots, \ell. \end{cases} \quad (4.38)$$

Then (4.36) and (4.37) can be written as

$$I^+ f(x) = \sum_{\alpha=1}^v \left\{ \left(\frac{\partial}{\partial \lambda} \right)^{n_\alpha-1} [a_\alpha(\lambda) \varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_\alpha} + \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{\lambda e_x^+(0, \lambda)}{e^+(\lambda)} F^+ \{ \varphi(x, \lambda) \varphi(f, \lambda) \} d\lambda, \quad (4.39)$$

$$I^- f(x) = - \sum_{\alpha=v+1}^{\ell} \left\{ \left(\frac{\partial}{\partial \lambda} \right)^{n_\alpha-1} [a_\alpha(\lambda) \varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_\alpha} + \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{\lambda e_x^-(0, \lambda)}{e^-(\lambda)} F^- \{ \varphi(x, \lambda) \varphi(f, \lambda) \} d\lambda. \quad (4.40)$$

We shall also use the following integral operators (see (4.7)):

$$Q^+ f(x) = \frac{1}{2\pi i} \int_{\Gamma_+} \frac{e_x^+(0, \lambda)}{e^+(\lambda)} \varphi(x, \lambda) \varphi(f, \lambda) d\lambda, \quad (4.41)$$

$$Q^- f(x) = \frac{1}{2\pi i} \int_{\Gamma_-} \frac{e_x^-(0, \lambda)}{e^-(\lambda)} \varphi(x, \lambda) \varphi(f, \lambda) d\lambda, \quad (4.42)$$

$$J^+ f(x) = \frac{1}{2\pi i} \sum_{\alpha=1}^v \sum_{\beta=0}^{n_\alpha-1} \left\{ \left(\frac{\partial}{\partial \lambda} \right)^\beta [\varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_\alpha} \times \int_{\Gamma_+} \frac{e_x^+(0, \lambda)}{e^+(\lambda)} F_{\alpha\beta}^+(\lambda) d\lambda + \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{e_x^+(0, \lambda)}{e^+(\lambda)} F^+ \{ \varphi(x, \lambda) \varphi(f, \lambda) \} d\lambda,$$

and

$$\begin{aligned}
 J^-f(x) &= \frac{1}{2\pi i} \sum_{\alpha=v+1}^{\ell} \sum_{\beta=0}^{n_{\alpha}-1} \left\{ \left(\frac{\partial}{\partial \lambda} \right)^{\beta} [\varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_{\alpha}} \\
 &\quad \times \int_{\Gamma_+} \frac{e_x^-(0, \lambda)}{e^-(\lambda)} F_{\alpha\beta}^-(\lambda) d\lambda \\
 &\quad + \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{e_x^-(0, \lambda)}{e^-(\lambda)} F^- \{ \varphi(x, \lambda) \varphi(f, \lambda) \} d\lambda.
 \end{aligned}$$

It is evident that,

$$Q^{\pm}f = J^{\pm}f,$$

for $f \in C_0^{\infty}(\mathbf{R}_+)$.

Similar to Lemma 4.4, we find

LEMMA 4.5. *For every $f \in H_{(n_0+1)}$ there exist a constant $c > 0$ such that*

$$\|J^{\pm}f\|_{-(n_0+1)} \leq c \|f\|_{(n_0+1)}.$$

It is evident that, for every $f \in H_{(n_0+1)}$

$$\begin{aligned}
 J^+f(x) &= \sum_{\alpha=1}^v \left\{ \left(\frac{\partial}{\partial \lambda} \right)^{n_{\alpha}-1} [b_{\alpha}(\lambda) \varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_{\alpha}} \\
 &\quad + \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{e_x^+(0, \lambda)}{e^+(\lambda)} F^+ \{ \varphi(x, \lambda) \varphi(f, \lambda) \} d\lambda, \quad (4.43)
 \end{aligned}$$

$$\begin{aligned}
 J^-f(x) &= - \sum_{\alpha=v+1}^{\ell} \left\{ \left(\frac{\partial}{\partial \lambda} \right)^{n_{\alpha}-1} [b_{\alpha}(\lambda) \varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_{\alpha}} \\
 &\quad + \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{e_x^-(0, \lambda)}{e^-(\lambda)} F^- \{ \varphi(x, \lambda) \varphi(f, \lambda) \} d\lambda, \quad (4.44)
 \end{aligned}$$

where

$$\begin{aligned}
 &\left\{ \left(\frac{d}{d\lambda} \right)^{n_{\alpha}-1-\beta} b_{\alpha}(\lambda) \right\}_{\lambda=\lambda_{\alpha}} \\
 &= \begin{cases} \frac{1}{2\pi i} \binom{n_{\alpha}-1}{\beta} \int_{\Gamma_+} \frac{e_x^+(0, \lambda)}{e^+(\lambda)} F_{\alpha\beta}^+(\lambda) d\lambda, & \alpha = 1, \dots, v \\ -\frac{1}{2\pi i} \binom{n_{\alpha}-1}{\beta} \int_{\Gamma_-} \frac{e_x^-(0, \lambda)}{e^-(\lambda)} F_{\alpha\beta}^-(\lambda) d\lambda, & \alpha = v+1, \dots, \ell. \end{cases} \quad (4.45)
 \end{aligned}$$

THEOREM 4.6. *Under the condition (1.4) the following two-fold spectral expansion in terms of the principal functions of L holds,*

$$\begin{aligned}
f(x) = & \sum_{i=1}^j \left\{ \left(\frac{\partial}{\partial \lambda} \right)^{m_i^+ - 1} [a_i^+(\lambda) \varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_i^+} \\
& + \sum_{i=1}^k \left\{ \left(\frac{\partial}{\partial \lambda} \right)^{m_i^- - 1} [a_i^-(\lambda) \varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_i^-} \\
& + \sum_{\alpha=1}^{\ell} \left\{ \left(\frac{\partial}{\partial \lambda} \right)^{n_{\alpha} - 1} [a_{\alpha}(\lambda) \varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_{\alpha}} \\
& + \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{\lambda e_x^+(0, \lambda)}{e^+(\lambda)} F^+ \{ \varphi(x, \lambda) \varphi(f, \lambda) \} d\lambda \\
& - \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{\lambda e_x^-(0, \lambda)}{e^-(\lambda)} F^- \{ \varphi(x, \lambda) \varphi(f, \lambda) \} d\lambda, \tag{4.46}
\end{aligned}$$

$$\begin{aligned}
0 = & \sum_{i=1}^j \left\{ \left(\frac{\partial}{\partial \lambda} \right)^{m_i^+ - 1} [b_i^+(\lambda) \varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_i^+} \\
& + \sum_{i=1}^k \left\{ \left(\frac{\partial}{\partial \lambda} \right)^{m_i^- - 1} [b_i^-(\lambda) \varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_i^-} \\
& + \sum_{\alpha=1}^{\ell} \left\{ \left(\frac{\partial}{\partial \lambda} \right)^{n_{\alpha} - 1} [b_{\alpha}(\lambda) \varphi(x, \lambda) \varphi(f, \lambda)] \right\}_{\lambda=\lambda_{\alpha}} \\
& + \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{e_x^+(0, \lambda)}{e^+(\lambda)} F^+ \{ \varphi(x, \lambda) \varphi(f, \lambda) \} d\lambda \\
& - \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{e_x^-(0, \lambda)}{e^-(\lambda)} F^- \{ \varphi(x, \lambda) \varphi(f, \lambda) \} d\lambda. \tag{4.47}
\end{aligned}$$

for every function $f \in H_{(n_0+1)}$. The integrals in (4.46) and (4.47) converge in the norm of $H_{-(n_0+1)}$ where a_i^{\pm} , b_i^{\pm} , F^+ , F^- , a_{α} , and b_{α} defined by (4.8), (4.9), (4.21), (4.22), (4.38), and (4.45), respectively.

Proof. We obtain (4.46) and (4.47) for $f \in C_0^{\infty}(\mathbf{R}_+) \subset H_{(n_0+1)}$, by use of (4.6), (4.7), (4.23), (4.24), and (4.39)–(4.44). The convergence of the integrals appearing in (4.46) and (4.47) in the norm of $H_{-(n_0+1)}$ has been given in Lemmas 4.4 and 4.5. Since $C_0^{\infty}(\mathbf{R}_+)$ is dense in $H_{(n_0+1)}$, the proof is completed. ■

Note. By (3.1), for every $f \in H_{(n_0+1)} \subset L^2(\mathbf{R}_+)$ the integrals

$$\varphi^{(m)}(f, \lambda_i^+) := \int_0^{\infty} f(x) \varphi^{(m)}(x, \lambda_i^+) dx, \quad m=0, 1, \dots, m_i^+ - 1, \quad i=1, \dots, j, \tag{4.48}$$

and

$$\varphi^{(m)}(f, \lambda_i^-) := \int_0^\infty f(x) \varphi^{(m)}(x, \lambda_i^-) dx, \quad m = 0, 1, \dots, m_i^- - 1, \quad i = 1, \dots, k, \quad (4.49)$$

converge. The differentiation in the first and second terms on the right-hand side of (4.46) and (4.47) is symbolic: that is, the rule for differentiating a product is first applied as if the derivatives of $\varphi(f, \lambda)$ at the points λ_i^+ and λ_i^- existed, and then these “derivatives” are replaced by the numbers $\varphi^{(m)}(f, \lambda_i^+)$ and $\varphi^{(m)}(f, \lambda_i^-)$ defined by (4.48) and (4.49), respectively. By Lemma 4.3 for every $f \in H_{(n_0+1)}$, the function $\varphi(f, \lambda)$ may be differentiated $(n_0 + 1)$ times with respect to λ ($\lambda \in \mathbf{R}$).

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REFERENCES

2. A. Degasperis, On the inverse problem for the Klein–Gordon s -wave equation, *J. Math. Phys.* **11** (1970), 551–567.
3. M. G. Gasymov, On the decomposition in a series of eigenfunctions for a non-selfadjoint boundary value problem of the solution of a differential equation with a singularity at zero point, *Soviet Math. Dokl.* **6** (1965), 1426–1429.
4. M. G. Gasymov, Expansion in terms of the solutions of a scattering theory problem for the non-selfadjoint Schrödinger equation, *Soviet Math. Dokl.* **9** (1968), 390–393.
5. M. G. Gasymov and F. G. Maksudov, The principal part of the resolvent of non-self-adjoint operators in neighbourhood of spectral singularities, *Funct. Anal. Appl.* **6** (1972), 185–192.
6. W. Greiner, “Relativistic Quantum Mechanics, Wave Equations,” Springer-Verlag, Berlin/New York, 1994.
7. S. V. Hruscev, Spectral singularities of dissipative Schrödinger operator with rapidly decreasing potential, *Indiana Univ. Math. J.* **33** (1984), 313–338.
8. M. Jaulent and C. Jean, The inverse s -wave scattering problem for a class of potentials depending on energy, *Comm. Math. Phys.* **28** (1972), 177–220.
9. M. Jaulent and C. Jean, The inverse problem for the one-dimensional Schrödinger equation with an energy-dependent potential I, II, *Ann. Inst. H. Poincaré Sec. A* **25** (1976), 105–118, 119–137.
10. M. V. Keldysh, On the completeness of the eigenfunctions of some classes of non-self-adjoint linear operators, *Soviet Math. Dokl.* **77** (1951), 11–14, *Russian Math. Survey* **26** (1971), 15–44.
11. R. R. D. Kemp, A singular boundary value problem for a non-selfadjoint differential operator, *Canad. J. Math.* **10** (1958), 447–462.
12. A. M. Krall, The adjoint of differential operator with integral boundary conditions, *Proc. Amer. Math. Soc.* **16** (1965), 738–742.

13. A. M. Krall, A nonhomogeneous eigenfunction expansion, *Trans. Amer. Math. Soc.* **117** (1965), 352–361.
14. A. M. Krall, Second order ordinary differential operators with general boundary conditions, *Duke J. Math.* **32** (1965), 617–625.
15. A. M. Krall, Nonhomogeneous differential operators, *Michigan Math. J.* **12** (1965), 247–265.
16. A. M. Krall, On non-selfadjoint ordinary differential operators of second order, *Soviet Math. Dokl.* **165** (1965), 1235–1237.
17. A. M. Krall, E. Bairamov, and O. Çakar, Spectrum and spectral singularities of quadratic pencil of Schrödinger operators with general boundary condition, *J. Differential Equations* **151** (1999), 252–267.
18. P. D. Lax and R. S. Phillips, “Scattering Theory,” Academic Press, San Diego, 1967.
19. V. E. Lyance, “A Differential Operator with Spectral Singularities, I, II,” Amer. Math. Soc. Transl., Ser. 2, Vol. 60, pp. 185–225, 227–283, Amer. Math. Soc., Providence, 1967.
20. F. G. Maksudov, Multiple expansion in the eigen and associated functions of a quadratic pencil of one-dimensional singular differential operators, “Spectral Theory of Operators,” pp. 125–162, Baku, 1977. [In Russian]
21. F. G. Maksudov and B. P. Allakhverdiev, Spectral analysis of a new class of non-self-adjoint operators with continuous and point spectrum, *Soviet Math. Dokl.* **30** (1984), 566–569.
22. F. G. Maksudov and G. Sh. Guseinov, On solution of the inverse scattering problem for a quadratic pencil of one-dimensional Schrödinger operators on the whole axis, *Soviet Math. Dokl.* **34** (1987), 34–38.
23. V. A. Marchenko, “Expansion in Eigenfunctions of Non-Selfadjoint Singular Second Order Differential Operators,” Amer. Math. Soc. Transl. Ser. 2, Vol. 25, pp. 77–130, Amer. Math. Soc., Providence, 1963.
24. B. Sz. Nagy and C. Foias, “Harmonic Analysis of Operators in Hilbert Space,” North-Holland, Amsterdam, 1970.
25. M. A. Naimark, Investigation of the spectrum and the expansion in eigenfunctions of a non-selfadjoint operator of second order on a semi-axis, Amer. Math. Soc. Transl. Ser. 2, Vol. 16, pp. 103–193, Amer. Math. Soc., Providence, 1960.
26. M. A. Naimark, “Linear Differential Operators I, II,” Ungar, New York, 1968.
27. B. S. Pavlov, The non-selfadjoint Schrödinger operator, *Topics in Math. Phys.* **1** (1967), 87–110.
28. B. S. Pavlov, On separation conditions for the spectral components of a dissipative operators, *Math. USSR Izvestiya* **9** (1975), 113–137.
29. J. T. Schwartz, Some non-selfadjoint operators, *Comm. Pure Appl. Math.* **13** (1960), 609–639.
30. E. C. Titchmarsh, “Eigenfunction Expansions Associated with Second Order Differential Equations,” Oxford Univ. Press, London, 1962.
31. O. A. Veliev, Spectral expansion of non-selfadjoint differential operators with periodic coefficients, *J. Differential Equations* **22** (1987), 1403–1408.